

Worskhop on Dynamics of Self-Gravitating Matter  
Institut Henri Poincaré, Paris

## Spin and Topo-spin

Topological aspects of Kerr-Newman spacetimes  
and their consequences for the Dirac equation

A. Shadi Tahvildar-Zadeh

Department of Mathematics  
Rutgers –New Brunswick

October 26, 2015

# Some Recent Results in GRQM

- ① A. S. Tahvildar-Zadeh, "On a zero-gravity limit of Kerr–Newman spacetimes and their electromagnetic fields," 22 pp. (2012) submitted to *Jour. Math. Phys.* (Nov. 2014) [[arXiv:1410.0416](#)].

# Some Recent Results in GRQM

- ① A. S. Tahvildar-Zadeh, "On a zero-gravity limit of Kerr–Newman spacetimes and their electromagnetic fields," 22 pp. (2012) submitted to *Jour. Math. Phys.* (Nov. 2014) [[arXiv:1410.0416](#)].
- ② M. K.-H. Kiessling & A. S. Tahvildar-Zadeh, "The Dirac point electron in zero-gravity Kerr–Newman spacetime," 47 pp., submitted to *Jour. Math. Phys.* (Nov 2014) [[arXiv:1410.0419](#)].

# Some Recent Results in GRQM

- ① A. S. Tahvildar-Zadeh, "On a zero-gravity limit of Kerr–Newman spacetimes and their electromagnetic fields," 22 pp. (2012) submitted to *Jour. Math. Phys.* (Nov. 2014) [[arXiv:1410.0416](#)].
- ② M. K.-H. Kiessling & A. S. Tahvildar-Zadeh, "The Dirac point electron in zero-gravity Kerr–Newman spacetime," 47 pp., submitted to *Jour. Math. Phys.* (Nov 2014) [[arXiv:1410.0419](#)].
- ③ M. K.-H. Kiessling & A. S. Tahvildar-Zadeh, "A novel quantum-mechanical interpretation of the Dirac equation," 47 pp., submitted to *Ann. of Phys.* (Nov 2014) [[arXiv:1411.2296](#)].

# Some Recent Results in GRQM

- ① A. S. Tahvildar-Zadeh, "On a zero-gravity limit of Kerr–Newman spacetimes and their electromagnetic fields," 22 pp. (2012) submitted to *Jour. Math. Phys.* (Nov. 2014) [[arXiv:1410.0416](#)].
- ② M. K.-H. Kiessling & A. S. Tahvildar-Zadeh, "The Dirac point electron in zero-gravity Kerr–Newman spacetime," 47 pp., submitted to *Jour. Math. Phys.* (Nov 2014) [[arXiv:1410.0419](#)].
- ③ M. K.-H. Kiessling & A. S. Tahvildar-Zadeh, "A novel quantum-mechanical interpretation of the Dirac equation," 47 pp., submitted to *Ann. of Phys.* (Nov 2014) [[arXiv:1411.2296](#)].
- ④ M. K. Balasubramanian, "Scalar fields and Spin-1/2 Fields on Mildly Singular Spacetimes," 71 pp., Ph.D. Dissertation, Rutgers –New Brunswick (May 2015).

# General-Relativistic Quantum Mechanics

- What do we mean by it? (not UT/GUT, not QG)

# General-Relativistic Quantum Mechanics

- What do we mean by it? (not UT/GUT, not QG)
- Are we even allowed to put those words together?

# General-Relativistic Quantum Mechanics

- What do we mean by it? (not UT/GUT, not QG)
- Are we even allowed to put those words together?
- *Macro vs. Micro*: Freeman Dyson's "Incompatible Worldviews"

# General-Relativistic Quantum Mechanics

- What do we mean by it? (not UT/GUT, not QG)
- Are we even allowed to put those words together?
- *Macro vs. Micro*: Freeman Dyson's "Incompatible Worldviews"

# General-Relativistic Quantum Mechanics

- What do we mean by it? (not UT/GUT, not QG)
- Are we even allowed to put those words together?
- *Macro vs. Micro: Freeman Dyson's "Incompatible Worldviews"*
- *"If this hypothesis were true, it would imply that theories of quantum gravity are untestable and **scientifically meaningless** ... and the search for a unified theory could turn out to be an illusion."*

# General-Relativistic Quantum Mechanics

- What do we mean by it? (not UT/GUT, not QG)
- Are we even allowed to put those words together?
- *Macro vs. Micro*: Freeman Dyson's "Incompatible Worldviews"
- *"If this hypothesis were true, it would imply that theories of quantum gravity are untestable and **scientifically meaningless** ... and the search for a unified theory could turn out to be an illusion."*
- Our question: Do these two theories have a common (mathematical) problem, that can perhaps be addressed by putting them together, like two pieces of the same puzzle?

# General-Relativistic Quantum Mechanics

- What do we mean by it? (not UT/GUT, not QG)
- Are we even allowed to put those words together?
- *Macro vs. Micro*: Freeman Dyson's "Incompatible Worldviews"
- *"If this hypothesis were true, it would imply that theories of quantum gravity are untestable and **scientifically meaningless** ... and the search for a unified theory could turn out to be an illusion."*
- Our question: Do these two theories have a common (mathematical) problem, that can perhaps be addressed by putting them together, like two pieces of the same puzzle?
- Fundamental objects that are highly singular in both gravity and electromagnetism: point mass and point charge.

# The Singularity Problem in GR

- The famous “Bridge” paper of Einstein & Rosen

# The Singularity Problem in GR

- The famous “Bridge” paper of Einstein & Rosen
- They glued two copies of exterior Schwarzschild together at their horizons, to obtain a spacetime free of singularities. But why?

# The Singularity Problem in GR

- The famous “Bridge” paper of Einstein & Rosen
- They glued two copies of exterior Schwarzschild together at their horizons, to obtain a spacetime free of singularities. But why?
- Title: “The Particle Problem in the General Theory of Relativity”, *Phys. Rev.* **48** (1935).

# The Singularity Problem in GR

- The famous “Bridge” paper of Einstein & Rosen
- They glued two copies of exterior Schwarzschild together at their horizons, to obtain a spacetime free of singularities. But why?
- Title: “The Particle Problem in the General Theory of Relativity”, *Phys. Rev.* **48** (1935).
- “To what extent can general relativity account for atomic structure of matter and for quantum effects?”

# The Singularity Problem in GR

- The famous “Bridge” paper of Einstein & Rosen
- They glued two copies of exterior Schwarzschild together at their horizons, to obtain a spacetime free of singularities. But why?
- Title: “The Particle Problem in the General Theory of Relativity”, *Phys. Rev.* **48** (1935).
- “To what extent can general relativity account for atomic structure of matter and for quantum effects?”
- What they wanted: solutions of Einstein’s equations that have particle-like features

# The Singularity Problem in GR

- The famous “Bridge” paper of Einstein & Rosen
- They glued two copies of exterior Schwarzschild together at their horizons, to obtain a spacetime free of singularities. But why?
- Title: “The Particle Problem in the General Theory of Relativity”, *Phys. Rev.* **48** (1935).
- “To what extent can general relativity account for atomic structure of matter and for quantum effects?”
- What they wanted: solutions of Einstein’s equations that have particle-like features
- Problem: They abhor singularities, so points are not allowed!

# The Singularity Problem in GR

- The famous “Bridge” paper of Einstein & Rosen
- They glued two copies of exterior Schwarzschild together at their horizons, to obtain a spacetime free of singularities. But why?
- Title: “The Particle Problem in the General Theory of Relativity”, *Phys. Rev.* **48** (1935).
- “To what extent can general relativity account for atomic structure of matter and for quantum effects?”
- What they wanted: solutions of Einstein’s equations that have particle-like features
- Problem: They abhor singularities, so points are not allowed!

# The Singularity Problem in GR

- The famous “Bridge” paper of Einstein & Rosen
- They glued two copies of exterior Schwarzschild together at their horizons, to obtain a spacetime free of singularities. But why?
- Title: “The Particle Problem in the General Theory of Relativity”, *Phys. Rev.* **48** (1935).
- “To what extent can general relativity account for atomic structure of matter and for quantum effects?”
- What they wanted: solutions of Einstein’s equations that have particle-like features
- Problem: They abhor singularities, so points are not allowed!
- *“Writers have occasionally noted the possibility that material particles might be considered as singularities of the field. This point of view, however, we cannot accept at all...Every field theory must adhere to the fundamental principle that singularities of the field are to be excluded.”*

# Particles as Bridges

- Einstein and Rosen come up with a solution that can represent a single particle, without having any singularity:

# Particles as Bridges

- Einstein and Rosen come up with a solution that can represent a single particle, without having any singularity:

# Particles as Bridges

- Einstein and Rosen come up with a solution that can represent a single particle, without having any singularity:
- *“These solutions involve the mathematical representation of physical space by a **space of two identical sheets, a [charged] particle being represented by a 'bridge' connecting these sheets.**”*

# Particles as Bridges

- Einstein and Rosen come up with a solution that can represent a single particle, without having any singularity:
- *“These solutions involve the mathematical representation of physical space by a **space of two identical sheets, a [charged] particle being represented by a 'bridge' connecting these sheets.**”*
- Two crazy ideas of E.&R.:

# Particles as Bridges

- Einstein and Rosen come up with a solution that can represent a single particle, without having any singularity:
- *“These solutions involve the mathematical representation of physical space by a **space of two identical sheets, a [charged] particle being represented by a 'bridge' connecting these sheets.**”*
- Two crazy ideas of E.&R.:
  - ① Particles can be represented within GR as bridges connecting two vacuums

# Particles as Bridges

- Einstein and Rosen come up with a solution that can represent a single particle, without having any singularity:
- *“These solutions involve the mathematical representation of physical space by a **space of two identical sheets, a [charged] particle being represented by a 'bridge' connecting these sheets.**”*
- Two crazy ideas of E.&R.:
  - 1 Particles can be represented within GR as bridges connecting two vacuums
  - 2 Space can be multi-sheeted (e.g. like a Riemann surface)

# Particles as Bridges

- Einstein and Rosen come up with a solution that can represent a single particle, without having any singularity:
- *“These solutions involve the mathematical representation of physical space by a **space of two identical sheets, a [charged] particle being represented by a 'bridge' connecting these sheets.**”*
- Two crazy ideas of E.&R.:
  - 1 Particles can be represented within GR as bridges connecting two vacuums
  - 2 Space can be multi-sheeted (e.g. like a Riemann surface)
- Who were those “writers” who liked singularities?

# Particles as Singularities

- Herman Weyl's Singularity Theory of Matter (1921):

# Particles as Singularities

- Herman Weyl's Singularity Theory of Matter (1921):

# Particles as Singularities

- Herman Weyl's Singularity Theory of Matter (1921):
- *"Matter now appears as a real singularity of the field (...) In the general theory of relativity the world can possess arbitrary (...) connectedness: nothing excludes the assumption that in its Analysis-Situs [i.e. topological] properties it behaves like a four dimensional Euclidean continuum, from which different tubes of infinite length in one dimension are cut off."*

# Particles as Singularities

- Herman Weyl's Singularity Theory of Matter (1921):
- *"Matter now appears as a real singularity of the field (...) In the general theory of relativity the world can possess arbitrary (...) connectedness: nothing excludes the assumption that in its Analysis-Situs [i.e. topological] properties it behaves like a four dimensional Euclidean continuum, from which different tubes of infinite length in one dimension are cut off."*
- Weyl's crazy ideas:

# Particles as Singularities

- Herman Weyl's Singularity Theory of Matter (1921):
- *"Matter now appears as a real singularity of the field (...) In the general theory of relativity the world can possess arbitrary (...) connectedness: nothing excludes the assumption that in its Analysis-Situs [i.e. topological] properties it behaves like a four dimensional Euclidean continuum, from which different tubes of infinite length in one dimension are cut off."*
- Weyl's crazy ideas:
  - 1 A particle *IS* a singularity of space

# Particles as Singularities

- Herman Weyl's Singularity Theory of Matter (1921):
- *"Matter now appears as a real singularity of the field (...) In the general theory of relativity the world can possess arbitrary (...) connectedness: nothing excludes the assumption that in its Analysis-Situs [i.e. topological] properties it behaves like a four dimensional Euclidean continuum, from which different tubes of infinite length in one dimension are cut off."*
- Weyl's crazy ideas:
  - ① A particle *IS* a singularity of space
  - ② These singularities need to be excised from spacetime: Matter *is* where space *isn't*.

# Particles as Singularities

- Herman Weyl's Singularity Theory of Matter (1921):
- *"Matter now appears as a real singularity of the field (...) In the general theory of relativity the world can possess arbitrary (...) connectedness: nothing excludes the assumption that in its Analysis-Situs [i.e. topological] properties it behaves like a four dimensional Euclidean continuum, from which different tubes of infinite length in one dimension are cut off."*
- Weyl's crazy ideas:
  - ① A particle *IS* a singularity of space
  - ② These singularities need to be excised from spacetime: Matter *is* where space *isn't*.
- Note: This also leads to spacetimes with non-trivial topology.

# Particles as Singularities

- Herman Weyl's Singularity Theory of Matter (1921):
- *"Matter now appears as a real singularity of the field (...) In the general theory of relativity the world can possess arbitrary (...) connectedness: nothing excludes the assumption that in its Analysis-Situs [i.e. topological] properties it behaves like a four dimensional Euclidean continuum, from which different tubes of infinite length in one dimension are cut off."*
- Weyl's crazy ideas:
  - ① A particle *IS* a singularity of space
  - ② These singularities need to be excised from spacetime: Matter *is* where space *isn't*.
- Note: This also leads to spacetimes with non-trivial topology.
- Ontology of Nothingness (*néant*-ology)

# Particles as Singularities

- Herman Weyl's Singularity Theory of Matter (1921):
- *"Matter now appears as a real singularity of the field (...) In the general theory of relativity the world can possess arbitrary (...) connectedness: nothing excludes the assumption that in its Analysis-Situs [i.e. topological] properties it behaves like a four dimensional Euclidean continuum, from which different tubes of infinite length in one dimension are cut off."*
- Weyl's crazy ideas:
  - ① A particle *IS* a singularity of space
  - ② These singularities need to be excised from spacetime: Matter *is* where space *isn't*.
- Note: This also leads to spacetimes with non-trivial topology.
- Ontology of Nothingness (*néant*-ology)
- PS: If you dislike a Weyl theory of matter, just wait five minutes!

# The Singularity Problem in QM

- Quantum law of motion for charged, point-like particles interacting with an electromagnetic field

# The Singularity Problem in QM

- Quantum law of motion for charged, point-like particles interacting with an electromagnetic field
- Electrodynamics:

# The Singularity Problem in QM

- Quantum law of motion for charged, point-like particles interacting with an electromagnetic field
- Electrodynamics:
  - ① Charges and currents are sources for the EM field (Maxwell)

# The Singularity Problem in QM

- Quantum law of motion for charged, point-like particles interacting with an electromagnetic field
- Electrodynamics:
  - 1 Charges and currents are sources for the EM field (Maxwell)
  - 2 EM Field act on charges and make the move (Lorentz)

# The Singularity Problem in QM

- Quantum law of motion for charged, point-like particles interacting with an electromagnetic field
- Electrodynamics:
  - ① Charges and currents are sources for the EM field (Maxwell)
  - ② EM Field act on charges and make the move (Lorentz)
- A classical problem in electrodynamics of point particles:

# The Singularity Problem in QM

- Quantum law of motion for charged, point-like particles interacting with an electromagnetic field
- Electrodynamics:
  - ① Charges and currents are sources for the EM field (Maxwell)
  - ② EM Field act on charges and make the move (Lorentz)
- A classical problem in electrodynamics of point particles:
  - ① EM field of a point charge is undefined at the location of the charge

# The Singularity Problem in QM

- Quantum law of motion for charged, point-like particles interacting with an electromagnetic field
- Electrodynamics:
  - ① Charges and currents are sources for the EM field (Maxwell)
  - ② EM Field act on charges and make the move (Lorentz)
- A classical problem in electrodynamics of point particles:
  - ① EM field of a point charge is undefined at the location of the charge
  - ② Self-Energy of a point charge is infinite

# The Singularity Problem in QM

- Quantum law of motion for charged, point-like particles interacting with an electromagnetic field
- Electrodynamics:
  - ① Charges and currents are sources for the EM field (Maxwell)
  - ② EM Field act on charges and make the move (Lorentz)
- A classical problem in electrodynamics of point particles:
  - ① EM field of a point charge is undefined at the location of the charge
  - ② Self-Energy of a point charge is infinite
- Consequences:

# The Singularity Problem in QM

- Quantum law of motion for charged, point-like particles interacting with an electromagnetic field
- Electrodynamics:
  - ① Charges and currents are sources for the EM field (Maxwell)
  - ② EM Field act on charges and make the move (Lorentz)
- A classical problem in electrodynamics of point particles:
  - ① EM field of a point charge is undefined at the location of the charge
  - ② Self-Energy of a point charge is infinite
- Consequences:
  - ① Maxwell-Lorentz electrodynamics is not well-defined for point charges

# The Singularity Problem in QM

- Quantum law of motion for charged, point-like particles interacting with an electromagnetic field
- Electrodynamics:
  - ① Charges and currents are sources for the EM field (Maxwell)
  - ② EM Field act on charges and make the move (Lorentz)
- A classical problem in electrodynamics of point particles:
  - ① EM field of a point charge is undefined at the location of the charge
  - ② Self-Energy of a point charge is infinite
- Consequences:
  - ① Maxwell-Lorentz electrodynamics is not well-defined for point charges
  - ② Electrovacuum spacetimes are highly singular

# The Singularity Problem in QM

- Quantum law of motion for charged, point-like particles interacting with an electromagnetic field
- Electrodynamics:
  - ① Charges and currents are sources for the EM field (Maxwell)
  - ② EM Field act on charges and make the move (Lorentz)
- A classical problem in electrodynamics of point particles:
  - ① EM field of a point charge is undefined at the location of the charge
  - ② Self-Energy of a point charge is infinite
- Consequences:
  - ① Maxwell-Lorentz electrodynamics is not well-defined for point charges
  - ② Electrovacuum spacetimes are highly singular
- These problems are NOT resolved (in a mathematically rigorous way) by currently known extensions of QM (QED, QFT, ...)

# Our line of inquiry

- Are there solutions of Einstein's equations, possibly multi-sheeted, with singularities that have particle-like features?

# Our line of inquiry

- Are there solutions of Einstein's equations, possibly multi-sheeted, with singularities that have particle-like features?
- Can one prescribe a quantum law of motion for those singularities in a way that is fully consistent with relativity?

# Our line of inquiry

- Are there solutions of Einstein's equations, possibly multi-sheeted, with singularities that have particle-like features?
- Can one prescribe a quantum law of motion for those singularities in a way that is fully consistent with relativity?
- Are the predictions of such a theory in agreement with physical experiments?

# (Non-relativistic) Quantum Mechanics

- Schrödinger 1926

# (Non-relativistic) Quantum Mechanics

- Schrödinger 1926
- The motion of microscopic systems is guided by a *wave function*

# (Non-relativistic) Quantum Mechanics

- Schrödinger 1926
- The motion of microscopic systems is guided by a *wave function*
- The wave function of a single particle is defined on the space of all possible positions of that particle (configuration space)

# (Non-relativistic) Quantum Mechanics

- Schrödinger 1926
- The motion of microscopic systems is guided by a *wave function*
- The wave function of a single particle is defined on the space of all possible positions of that particle (configuration space)
- The wave function of a system of  $N$  particles is defined on the space of all possible positions of those  $N$  particles, i.e. on  $\mathbb{R}^{3N}$

# (Non-relativistic) Quantum Mechanics

- Schrödinger 1926
- The motion of microscopic systems is guided by a *wave function*
- The wave function of a single particle is defined on the space of all possible positions of that particle (configuration space)
- The wave function of a system of  $N$  particles is defined on the space of all possible positions of those  $N$  particles, i.e. on  $\mathbb{R}^{3N}$
- The wave function satisfies an evolution equation in a Hilbert space  $\mathcal{V}$

# (Non-relativistic) Quantum Mechanics

- Schrödinger 1926
- The motion of microscopic systems is guided by a *wave function*
- The wave function of a single particle is defined on the space of all possible positions of that particle (configuration space)
- The wave function of a system of  $N$  particles is defined on the space of all possible positions of those  $N$  particles, i.e. on  $\mathbb{R}^{3N}$
- The wave function satisfies an evolution equation in a Hilbert space  $\mathcal{V}$

$$i\frac{\partial}{\partial t}\Psi = H\Psi$$

# (Non-relativistic) Quantum Mechanics

- Schrödinger 1926
- The motion of microscopic systems is guided by a *wave function*
- The wave function of a single particle is defined on the space of all possible positions of that particle (configuration space)
- The wave function of a system of  $N$  particles is defined on the space of all possible positions of those  $N$  particles, i.e. on  $\mathbb{R}^{3N}$
- The wave function satisfies an evolution equation in a Hilbert space  $\mathcal{V}$

$$i\frac{\partial}{\partial t}\Psi = H\Psi$$

- $\Psi = \Psi(t, x), \Psi(t, \cdot) \in \mathcal{V}, H: \mathcal{V} \rightarrow \mathcal{V}$

# (Non-relativistic) Quantum Mechanics

- Schrödinger 1926
- The motion of microscopic systems is guided by a *wave function*
- The wave function of a single particle is defined on the space of all possible positions of that particle (configuration space)
- The wave function of a system of  $N$  particles is defined on the space of all possible positions of those  $N$  particles, i.e. on  $\mathbb{R}^{3N}$
- The wave function satisfies an evolution equation in a Hilbert space  $\mathcal{V}$

$$i\frac{\partial}{\partial t}\Psi = H\Psi$$

- $\Psi = \Psi(t, x)$ ,  $\Psi(t, \cdot) \in \mathcal{V}$ ,  $H: \mathcal{V} \rightarrow \mathcal{V}$
- If  $\Psi(t, x) = e^{-iEt}\psi(x)$  then  $H\psi = E\psi$

# (Non-relativistic) Quantum Mechanics

- Schrödinger 1926
- The motion of microscopic systems is guided by a *wave function*
- The wave function of a single particle is defined on the space of all possible positions of that particle (configuration space)
- The wave function of a system of  $N$  particles is defined on the space of all possible positions of those  $N$  particles, i.e. on  $\mathbb{R}^{3N}$
- The wave function satisfies an evolution equation in a Hilbert space  $\mathcal{V}$

$$i\frac{\partial}{\partial t}\Psi = H\Psi$$

- $\Psi = \Psi(t, x)$ ,  $\Psi(t, \cdot) \in \mathcal{V}$ ,  $H: \mathcal{V} \rightarrow \mathcal{V}$
- If  $\Psi(t, x) = e^{-iEt}\psi(x)$  then  $H\psi = E\psi$
- Eigenvalues and Eigenfunctions: Physics becomes Linear Algebra!

# Spectrum of Hydrogen

- What do you see when you heat up a gaseous element until it glows, and look at it through a prism?

# Spectrum of Hydrogen

- What do you see when you heat up a gaseous element until it glows, and look at it through a prism?
- Spectral lines.

# Spectrum of Hydrogen

- What do you see when you heat up a gaseous element until it glows, and look at it through a prism?
- Spectral lines.
- Why not a rainbow?

# Spectrum of Hydrogen

- What do you see when you heat up a gaseous element until it glows, and look at it through a prism?
- Spectral lines.
- Why not a rainbow?
- Bohr and the quantum revolution

# Spectrum of Hydrogen

- What do you see when you heat up a gaseous element until it glows, and look at it through a prism?
- Spectral lines.
- Why not a rainbow?
- Bohr and the quantum revolution
- Orbitals are Energy eigenstates!

# Spectrum of Hydrogen

- What do you see when you heat up a gaseous element until it glows, and look at it through a prism?
- Spectral lines.
- Why not a rainbow?
- Bohr and the quantum revolution
- Orbitals are Energy eigenstates!
- Schrödinger's equation reproduces Bohr's spectrum!

# Special-relativistic Quantum Mechanics

- Schrödinger's equation is non-relativistic

# Special-relativistic Quantum Mechanics

- Schrödinger's equation is non-relativistic
- Sommerfeld's relativistic corrections to Bohr's spectrum

# Special-relativistic Quantum Mechanics

- Schrödinger's equation is non-relativistic
- Sommerfeld's relativistic corrections to Bohr's spectrum
- Dirac 1928: New equation for the wave function

# Special-relativistic Quantum Mechanics

- Schrödinger's equation is non-relativistic
- Sommerfeld's relativistic corrections to Bohr's spectrum
- Dirac 1928: New equation for the wave function
- Dirac's equation is Lorentz-invariant

# Special-relativistic Quantum Mechanics

- Schrödinger's equation is non-relativistic
- Sommerfeld's relativistic corrections to Bohr's spectrum
- Dirac 1928: New equation for the wave function
- Dirac's equation is Lorentz-invariant
- Dirac's spectral lines came out to be exactly the same as Sommerfeld's corrected version of Bohr's.

# Special-relativistic Quantum Mechanics

- Schrödinger's equation is non-relativistic
- Sommerfeld's relativistic corrections to Bohr's spectrum
- Dirac 1928: New equation for the wave function
- Dirac's equation is Lorentz-invariant
- Dirac's spectral lines came out to be exactly the same as Sommerfeld's corrected version of Bohr's.
- Schrödinger and Dirac shared a Nobel Prize!

# Special-Relativistic Hydrogen

- Dirac's Equation for Electron in Proton's electrostatic Field, in Hamiltonian form:  $i\partial_t\Psi = \mathbf{H}\Psi$

# Special-Relativistic Hydrogen

- Dirac's Equation for Electron in Proton's electrostatic Field, in Hamiltonian form:  $i\partial_t\Psi = \mathbf{H}\Psi$
- Eigen-functions:  $\Psi(t, \mathbf{x}) = e^{-iEt}\psi(\mathbf{x}) \implies \mathbf{H}\psi = E\psi$

# Special-Relativistic Hydrogen

- Dirac's Equation for Electron in Proton's electrostatic Field, in Hamiltonian form:  $i\partial_t\Psi = \mathbf{H}\Psi$
- Eigen-functions:  $\Psi(t, \mathbf{x}) = e^{-iEt}\psi(\mathbf{x}) \implies \mathbf{H}\psi = E\psi$
- $\sigma_{disc}(H_{Dirac}) = \left\{ \frac{m}{\sqrt{1 + \frac{e^4}{(n-|k| + \sqrt{k^2 - e^4})^2}}} \right\}$

# Special-Relativistic Hydrogen

- Dirac's Equation for Electron in Proton's electrostatic Field, in Hamiltonian form:  $i\partial_t\Psi = \mathbf{H}\Psi$
- Eigen-functions:  $\Psi(t, \mathbf{x}) = e^{-iEt}\psi(\mathbf{x}) \implies \mathbf{H}\psi = E\psi$
- $\sigma_{disc}(H_{Dirac}) = \left\{ \frac{m}{\sqrt{1 + \frac{e^4}{(n-|k| + \sqrt{k^2 - e^4})^2}}} \right\}$
- $n = 1, 2, 3, \dots, \quad k = -n, -n+1, \dots, -1, 1, \dots, n-1$

# Special-Relativistic Hydrogen

- Dirac's Equation for Electron in Proton's electrostatic Field, in Hamiltonian form:  $i\partial_t\Psi = \mathbf{H}\Psi$
- Eigen-functions:  $\Psi(t, \mathbf{x}) = e^{-iEt}\psi(\mathbf{x}) \implies \mathbf{H}\psi = E\psi$
- $\sigma_{disc}(H_{Dirac}) = \left\{ \frac{m}{\sqrt{1 + \frac{e^4}{(n-|k| + \sqrt{k^2 - e^4})^2}}} \right\}$
- $n = 1, 2, 3, \dots, \quad k = -n, -n+1, \dots, -1, 1, \dots, n-1$
- Spectroscopy.  $s, p, d, f, g, \dots$  Orbitals. Degenerate and non-degenerate states. Hyperfine splitting. Lamb shift. QED. More Nobel prizes!

# Special-Relativistic Hydrogen

- Dirac's Equation for Electron in Proton's electrostatic Field, in Hamiltonian form:  $i\partial_t\psi = \mathbf{H}\psi$
- Eigen-functions:  $\psi(t, \mathbf{x}) = e^{-iEt}\psi(\mathbf{x}) \implies \mathbf{H}\psi = E\psi$
- $\sigma_{disc}(H_{Dirac}) = \left\{ \frac{m}{\sqrt{1 + \frac{e^4}{(n-|k| + \sqrt{k^2 - e^4})^2}}} \right\}$
- $n = 1, 2, 3, \dots, \quad k = -n, -n+1, \dots, -1, 1, \dots, n-1$
- Spectroscopy.  $s, p, d, f, g, \dots$  Orbitals. Degenerate and non-degenerate states. Hyperfine splitting. Lamb shift. QED. More Nobel prizes!
- $\sigma_{cont}(H_{Dirac}) = (-\infty, -m] \cup [m, \infty)$

# Special-Relativistic Hydrogen

- Dirac's Equation for Electron in Proton's electrostatic Field, in Hamiltonian form:  $i\partial_t\Psi = \mathbf{H}\Psi$
- Eigen-functions:  $\Psi(t, \mathbf{x}) = e^{-iEt}\psi(\mathbf{x}) \implies \mathbf{H}\psi = E\psi$
- $\sigma_{disc}(H_{Dirac}) = \left\{ \frac{m}{\sqrt{1 + \frac{e^4}{(n-|k| + \sqrt{k^2 - e^4})^2}}} \right\}$
- $n = 1, 2, 3, \dots, \quad k = -n, -n+1, \dots, -1, 1, \dots, n-1$
- Spectroscopy.  $s, p, d, f, g, \dots$  Orbitals. Degenerate and non-degenerate states. Hyperfine splitting. Lamb shift. QED. More Nobel prizes!
- $\sigma_{cont}(H_{Dirac}) = (-\infty, -m] \cup [m, \infty)$
- Negative Continuum: Dirac's Sea, Hole Theory, Positron

# General-relativistic Hydrogen?

- Einstein's dream

# General-relativistic Hydrogen?

- Einstein's dream
- Gravity is much weaker than electromagnetism, so adding gravitational effects to Dirac's equation should shift spectral lines by a tiny amount, if at all

# General-relativistic Hydrogen?

- Einstein's dream
- Gravity is much weaker than electromagnetism, so adding gravitational effects to Dirac's equation should shift spectral lines by a tiny amount, if at all
- Nasty Surprise: general-relativistic Dirac Hamiltonian has NO discrete spectrum. (Belgiorno et al, Finster et al.)

# General-relativistic Hydrogen?

- Einstein's dream
- Gravity is much weaker than electromagnetism, so adding gravitational effects to Dirac's equation should shift spectral lines by a tiny amount, if at all
- Nasty Surprise: general-relativistic Dirac Hamiltonian has NO discrete spectrum. (Belgiorno et al, Finster et al.)
- Pathologies of well-known solutions to Einstein's equations

# General-relativistic Hydrogen?

- Einstein's dream
- Gravity is much weaker than electromagnetism, so adding gravitational effects to Dirac's equation should shift spectral lines by a tiny amount, if at all
- Nasty Surprise: general-relativistic Dirac Hamiltonian has NO discrete spectrum. (Belgiorno et al, Finster et al.)
- Pathologies of well-known solutions to Einstein's equations
- Likely culprit: Infinite self-energy of point charges (linear electromagnetics) causes the spacetime to be infinitely curved close to the charge.

# General-relativistic Hydrogen?

- Einstein's dream
- Gravity is much weaker than electromagnetism, so adding gravitational effects to Dirac's equation should shift spectral lines by a tiny amount, if at all
- Nasty Surprise: general-relativistic Dirac Hamiltonian has NO discrete spectrum. (Belgiorno et al, Finster et al.)
- Pathologies of well-known solutions to Einstein's equations
- Likely culprit: Infinite self-energy of point charges (linear electromagnetics) causes the spacetime to be infinitely curved close to the charge.
- Two remedies:

# General-relativistic Hydrogen?

- Einstein's dream
- Gravity is much weaker than electromagnetism, so adding gravitational effects to Dirac's equation should shift spectral lines by a tiny amount, if at all
- Nasty Surprise: general-relativistic Dirac Hamiltonian has NO discrete spectrum. (Belgiorno et al, Finster et al.)
- Pathologies of well-known solutions to Einstein's equations
- Likely culprit: Infinite self-energy of point charges (linear electromagnetics) causes the spacetime to be infinitely curved close to the charge.
- Two remedies:
  - 1 Nonlinear electromagnetics can give rise to *finite* self-energies, and thus milder singularities for the spacetime

# General-relativistic Hydrogen?

- Einstein's dream
- Gravity is much weaker than electromagnetism, so adding gravitational effects to Dirac's equation should shift spectral lines by a tiny amount, if at all
- Nasty Surprise: general-relativistic Dirac Hamiltonian has NO discrete spectrum. (Belgiorno et al, Finster et al.)
- Pathologies of well-known solutions to Einstein's equations
- Likely culprit: Infinite self-energy of point charges (linear electromagnetics) causes the spacetime to be infinitely curved close to the charge.
- Two remedies:
  - 1 Nonlinear electromagnetics can give rise to *finite* self-energies, and thus milder singularities for the spacetime
  - 2 Zero-gravity limit: working with topological remnants of gravity

# Nonlinear Electromagnetics

- Maxwell's equations involve four fields: **E**, **B**, **D**, **H**.

# Nonlinear Electromagnetics

- Maxwell's equations involve four fields:  $\mathbf{E}$ ,  $\mathbf{B}$ ,  $\mathbf{D}$ ,  $\mathbf{H}$ .
- Need constitutive relations to close:  $\mathbf{E} = \mathbf{E}(\mathbf{D}, \mathbf{B})$ ,  $\mathbf{H} = \mathbf{H}(\mathbf{D}, \mathbf{B})$ .

# Nonlinear Electromagnetics

- Maxwell's equations involve four fields:  $\mathbf{E}$ ,  $\mathbf{B}$ ,  $\mathbf{D}$ ,  $\mathbf{H}$ .
- Need constitutive relations to close:  $\mathbf{E} = \mathbf{E}(\mathbf{D}, \mathbf{B})$ ,  $\mathbf{H} = \mathbf{H}(\mathbf{D}, \mathbf{B})$ .
- Maxwell's vacuum law:  $\mathbf{E} = \mathbf{D}$ ,  $\mathbf{H} = \mathbf{B}$ . Linear electromagnetics

# Nonlinear Electromagnetics

- Maxwell's equations involve four fields:  $\mathbf{E}$ ,  $\mathbf{B}$ ,  $\mathbf{D}$ ,  $\mathbf{H}$ .
- Need constitutive relations to close:  $\mathbf{E} = \mathbf{E}(\mathbf{D}, \mathbf{B})$ ,  $\mathbf{H} = \mathbf{H}(\mathbf{D}, \mathbf{B})$ .
- Maxwell's vacuum law:  $\mathbf{E} = \mathbf{D}$ ,  $\mathbf{H} = \mathbf{B}$ . Linear electromagnetics
- Coulomb potential and the electrostatic energy of a point charge

# Nonlinear Electromagnetics

- Maxwell's equations involve four fields:  $\mathbf{E}$ ,  $\mathbf{B}$ ,  $\mathbf{D}$ ,  $\mathbf{H}$ .
- Need constitutive relations to close:  $\mathbf{E} = \mathbf{E}(\mathbf{D}, \mathbf{B})$ ,  $\mathbf{H} = \mathbf{H}(\mathbf{D}, \mathbf{B})$ .
- Maxwell's vacuum law:  $\mathbf{E} = \mathbf{D}$ ,  $\mathbf{H} = \mathbf{B}$ . Linear electromagnetics
- Coulomb potential and the electrostatic energy of a point charge
- Born's idea:  $\mathbf{E} = \mathbf{D} / \sqrt{1 + |\mathbf{D}|^2}$ . Born-Infeld's nonlinear electromagnetics. Finite self-energy

# Nonlinear Electromagnetics

- Maxwell's equations involve four fields:  $\mathbf{E}$ ,  $\mathbf{B}$ ,  $\mathbf{D}$ ,  $\mathbf{H}$ .
- Need constitutive relations to close:  $\mathbf{E} = \mathbf{E}(\mathbf{D}, \mathbf{B})$ ,  $\mathbf{H} = \mathbf{H}(\mathbf{D}, \mathbf{B})$ .
- Maxwell's vacuum law:  $\mathbf{E} = \mathbf{D}$ ,  $\mathbf{H} = \mathbf{B}$ . Linear electromagnetics
- Coulomb potential and the electrostatic energy of a point charge
- Born's idea:  $\mathbf{E} = \mathbf{D} / \sqrt{1 + |\mathbf{D}|^2}$ . Born-Infeld's nonlinear electromagnetics. Finite self-energy
- Einstein-Maxwell-Born-Infeld equations. The Hoffmann spacetime. Conical singularities.

# Nonlinear Electromagnetics

- Maxwell's equations involve four fields:  $\mathbf{E}$ ,  $\mathbf{B}$ ,  $\mathbf{D}$ ,  $\mathbf{H}$ .
- Need constitutive relations to close:  $\mathbf{E} = \mathbf{E}(\mathbf{D}, \mathbf{B})$ ,  $\mathbf{H} = \mathbf{H}(\mathbf{D}, \mathbf{B})$ .
- Maxwell's vacuum law:  $\mathbf{E} = \mathbf{D}$ ,  $\mathbf{H} = \mathbf{B}$ . Linear electromagnetics
- Coulomb potential and the electrostatic energy of a point charge
- Born's idea:  $\mathbf{E} = \mathbf{D} / \sqrt{1 + |\mathbf{D}|^2}$ . Born-Infeld's nonlinear electromagnetics. Finite self-energy
- Einstein-Maxwell-Born-Infeld equations. The Hoffmann spacetime. Conical singularities.
- Dirac equation on Hoffmann spacetime: Self-adjointness of the Hamiltonian, and the existence of discrete spectrum (M.K. Balasubramanian, 2015)

# Nonlinear Electromagnetics

- Maxwell's equations involve four fields:  $\mathbf{E}$ ,  $\mathbf{B}$ ,  $\mathbf{D}$ ,  $\mathbf{H}$ .
- Need constitutive relations to close:  $\mathbf{E} = \mathbf{E}(\mathbf{D}, \mathbf{B})$ ,  $\mathbf{H} = \mathbf{H}(\mathbf{D}, \mathbf{B})$ .
- Maxwell's vacuum law:  $\mathbf{E} = \mathbf{D}$ ,  $\mathbf{H} = \mathbf{B}$ . Linear electromagnetics
- Coulomb potential and the electrostatic energy of a point charge
- Born's idea:  $\mathbf{E} = \mathbf{D} / \sqrt{1 + |\mathbf{D}|^2}$ . Born-Infeld's nonlinear electromagnetics. Finite self-energy
- Einstein-Maxwell-Born-Infeld equations. The Hoffmann spacetime. Conical singularities.
- Dirac equation on Hoffmann spacetime: Self-adjointness of the Hamiltonian, and the existence of discrete spectrum (M.K. Balasubramanian, 2015)
- The Great Challenge: going beyond spherical symmetry for Einstein-Maxwell-Born-Infeld.

# General-relativistic Zero-Gravity Quantum Mechanics

- Two sources of  $G$ -dependence for Einstein geometries

# General-relativistic Zero-Gravity Quantum Mechanics

- Two sources of  $G$ -dependence for Einstein geometries
- The coupling constant  $G$

# General-relativistic Zero-Gravity Quantum Mechanics

- Two sources of  $G$ -dependence for Einstein geometries
- The coupling constant  $G$
- The dimension conversion constant  $G$

# General-relativistic Zero-Gravity Quantum Mechanics

- Two sources of  $G$ -dependence for Einstein geometries
- The coupling constant  $G$
- The dimension conversion constant  $G$
- Example: Reissner-Nordström, in spherical coordinates

$$g_{RN} = \text{diag}(-f(r), \frac{1}{f(r)}, r^2, r^2 \sin^2 \theta)$$

# General-relativistic Zero-Gravity Quantum Mechanics

- Two sources of  $G$ -dependence for Einstein geometries
- The coupling constant  $G$
- The dimension conversion constant  $G$
- Example: Reissner-Nordström, in spherical coordinates

$$g_{RN} = \text{diag}(-f(r), \frac{1}{f(r)}, r^2, r^2 \sin^2 \theta)$$

- $f(r) = 1 - \frac{2GM}{c^2 r} + \frac{GQ^2}{c^4 r^2}$

# General-relativistic Zero-Gravity Quantum Mechanics

- Two sources of  $G$ -dependence for Einstein geometries
- The coupling constant  $G$
- The dimension conversion constant  $G$
- Example: Reissner-Nordström, in spherical coordinates  

$$g_{RN} = \text{diag}(-f(r), \frac{1}{f(r)}, r^2, r^2 \sin^2 \theta)$$
- $f(r) = 1 - \frac{2GM}{c^2 r} + \frac{GQ^2}{c^4 r^2}$
- Take the  $G \rightarrow 0$  limit, recover Minkowski space

# General-relativistic Zero-Gravity Quantum Mechanics

- Two sources of  $G$ -dependence for Einstein geometries
- The coupling constant  $G$
- The dimension conversion constant  $G$
- Example: Reissner-Nordström, in spherical coordinates  

$$g_{RN} = \text{diag}(-f(r), \frac{1}{f(r)}, r^2, r^2 \sin^2 \theta)$$
- $f(r) = 1 - \frac{2GM}{c^2 r} + \frac{GQ^2}{c^4 r^2}$
- Take the  $G \rightarrow 0$  limit, recover Minkowski space
- There are other famous solutions of Einstein's Eqs, where the zero- $G$  limit does NOT give you back Minkowski space

# General-relativistic Zero-Gravity Quantum Mechanics

- Two sources of  $G$ -dependence for Einstein geometries
- The coupling constant  $G$
- The dimension conversion constant  $G$
- Example: Reissner-Nordström, in spherical coordinates  

$$g_{RN} = \text{diag}(-f(r), \frac{1}{f(r)}, r^2, r^2 \sin^2 \theta)$$
- $f(r) = 1 - \frac{2GM}{c^2 r} + \frac{GQ^2}{c^4 r^2}$
- Take the  $G \rightarrow 0$  limit, recover Minkowski space
- There are other famous solutions of Einstein's Eqs, where the zero- $G$  limit does NOT give you back Minkowski space
- Nontrivial geometry goes away, nontrivial topology remains

# General-relativistic Zero-Gravity Quantum Mechanics

- Two sources of  $G$ -dependence for Einstein geometries
- The coupling constant  $G$
- The dimension conversion constant  $G$
- Example: Reissner-Nordström, in spherical coordinates  

$$g_{RN} = \text{diag}(-f(r), \frac{1}{f(r)}, r^2, r^2 \sin^2 \theta)$$
- $f(r) = 1 - \frac{2GM}{c^2 r} + \frac{GQ^2}{c^4 r^2}$
- Take the  $G \rightarrow 0$  limit, recover Minkowski space
- There are other famous solutions of Einstein's Eqs, where the zero- $G$  limit does NOT give you back Minkowski space
- Nontrivial geometry goes away, nontrivial topology remains
- The Kerr-Newman solution (1965), and its Maximal Analytical Extension (1968)

# General-relativistic Zero-Gravity Quantum Mechanics

- Two sources of  $G$ -dependence for Einstein geometries
- The coupling constant  $G$
- The dimension conversion constant  $G$
- Example: Reissner-Nordström, in spherical coordinates  

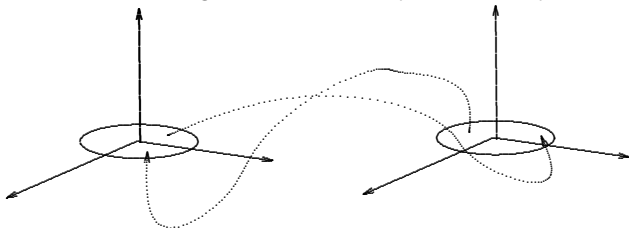
$$g_{RN} = \text{diag}(-f(r), \frac{1}{f(r)}, r^2, r^2 \sin^2 \theta)$$
- $f(r) = 1 - \frac{2GM}{c^2 r} + \frac{GQ^2}{c^4 r^2}$
- Take the  $G \rightarrow 0$  limit, recover Minkowski space
- There are other famous solutions of Einstein's Eqs, where the zero- $G$  limit does NOT give you back Minkowski space
- Nontrivial geometry goes away, nontrivial topology remains
- The Kerr-Newman solution (1965), and its Maximal Analytical Extension (1968)
- Appell-Sommerfeld: The zero-gravity limit of the maximal-analytically-extended Kerr-Newman solution

# Appell and Sommerfeld

- Sommerfeld (1896) Generalization of Riemann surfaces to 3-D

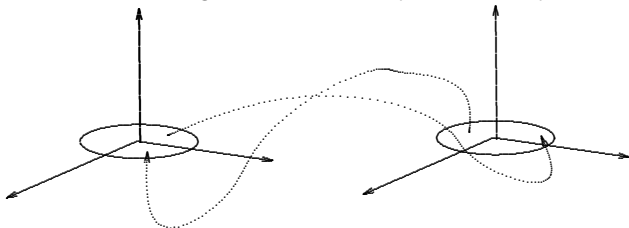
# Appell and Sommerfeld

- Sommerfeld (1896) Generalization of Riemann surfaces to 3-D
- Two copies of  $\mathbb{R}^3$  cross-glued at a disk (of radius  $a$ ):



# Appell and Sommerfeld

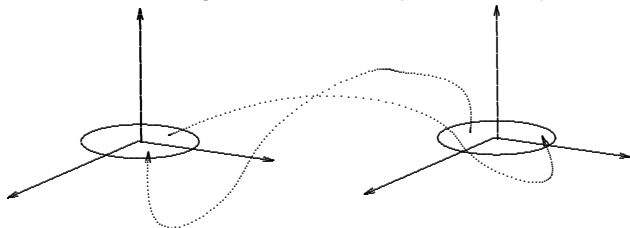
- Sommerfeld (1896) Generalization of Riemann surfaces to 3-D
- Two copies of  $\mathbb{R}^3$  cross-glued at a disk (of radius  $a$ ):



- Let's watch a video!

# Appell and Sommerfeld

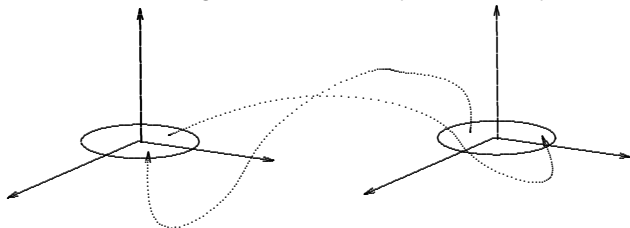
- Sommerfeld (1896) Generalization of Riemann surfaces to 3-D
- Two copies of  $\mathbb{R}^3$  cross-glued at a disk (of radius  $a$ ):



- Let's watch a video!
- There is a coordinate system that covers this manifold in a single chart: Oblate Spheroidal Coordinates  $(r, \theta, \varphi)$

# Appell and Sommerfeld

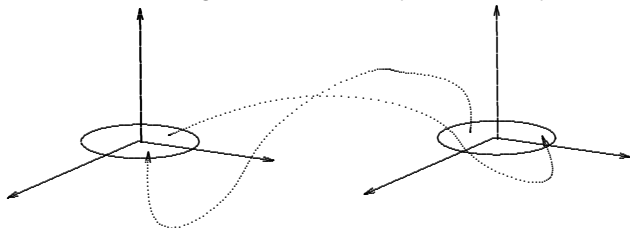
- Sommerfeld (1896) Generalization of Riemann surfaces to 3-D
- Two copies of  $\mathbb{R}^3$  cross-glued at a disk (of radius  $a$ ):



- Let's watch a video!
- There is a coordinate system that covers this manifold in a single chart: Oblate Spheroidal Coordinates  $(r, \theta, \varphi)$
- Add a time dimension to get a static spacetime.

# Appell and Sommerfeld

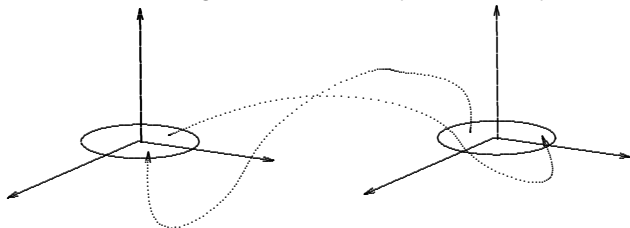
- Sommerfeld (1896) Generalization of Riemann surfaces to 3-D
- Two copies of  $\mathbb{R}^3$  cross-glued at a disk (of radius  $a$ ):



- Let's watch a video!
- There is a coordinate system that covers this manifold in a single chart: Oblate Spheroidal Coordinates  $(r, \theta, \varphi)$
- Add a time dimension to get a static spacetime.
- It is flat (away from the ring)  $\implies$  it is a solution of E.V.E.

# Appell and Sommerfeld

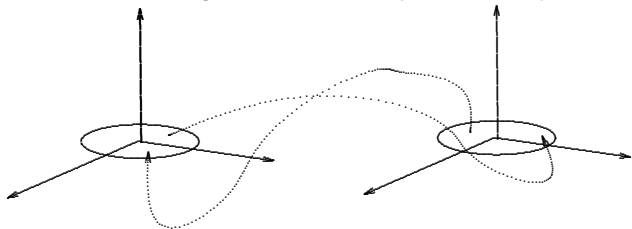
- Sommerfeld (1896) Generalization of Riemann surfaces to 3-D
- Two copies of  $\mathbb{R}^3$  cross-glued at a disk (of radius  $a$ ):



- Let's watch a video!
- There is a coordinate system that covers this manifold in a single chart: Oblate Spheroidal Coordinates  $(r, \theta, \varphi)$
- Add a time dimension to get a static spacetime.
- It is flat (away from the ring)  $\implies$  it is a solution of E.V.E.
- Points on the ring are conical singularities for the metric

# Appell and Sommerfeld

- Sommerfeld (1896) Generalization of Riemann surfaces to 3-D
- Two copies of  $\mathbb{R}^3$  cross-glued at a disk (of radius  $a$ ):



- Let's watch a video!
- There is a coordinate system that covers this manifold in a single chart: Oblate Spheroidal Coordinates  $(r, \theta, \varphi)$
- Add a time dimension to get a static spacetime.
- It is flat (away from the ring)  $\implies$  it is a solution of E.V.E.
- Points on the ring are conical singularities for the metric
- zGKN = This spacetime + EM fields on it (P. Appell, 1888)

# The Zero- $G$ Limit comes to Dirac's Rescue

- Dirac's equation for an electron in the zero- $G$  Kerr–Newman (zGKN) spacetime

# The Zero- $G$ Limit comes to Dirac's Rescue

- Dirac's equation for an electron in the zero- $G$  Kerr–Newman (zGKN) spacetime
- The zero- $G$  limit of a positively charged Reissner- Nordström spacetime yields Minkowski spacetime with a positive Coulomb singularity — **Nothing New!**: *Back to Special-Relativistic Hydrogen*

# The Zero- $G$ Limit comes to Dirac's Rescue

- Dirac's equation for an electron in the zero- $G$  Kerr–Newman (zGKN) spacetime
- The zero- $G$  limit of a positively charged Reissner- Nordström spacetime yields Minkowski spacetime with a positive Coulomb singularity — **Nothing New!**: *Back to Special-Relativistic Hydrogen*
- The zero- $G$  limit of the maximally extended Kerr-Newman spacetime yields a **flat**, but **topologically non-trivial**, **multi-sheeted** electromagnetic spacetime.

This does yield **Something New!** .... and Interesting:  
*Zero- $G$  Kerr-Newman Hydrogen exists!*

# Some of our results on zero- $G$ Hydrogen spectrum

- The Dirac Hamiltonian on zero- $G$  Kerr Newman spacetime is essentially self-adjoint

# Some of our results on zero- $G$ Hydrogen spectrum

- The Dirac Hamiltonian on zero- $G$  Kerr Newman spacetime is **essentially self-adjoint**
- The spectrum is **symmetric** about zero.

# Some of our results on zero- $G$ Hydrogen spectrum

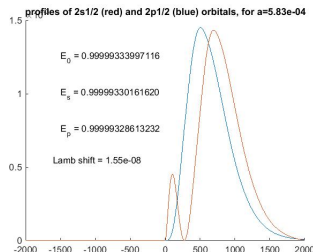
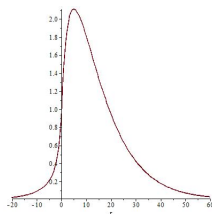
- The Dirac Hamiltonian on zero- $G$  Kerr Newman spacetime is **essentially self-adjoint**
- The spectrum is **symmetric** about zero.
- Discrete spectrum is **non-empty** under some smallness conditions.

# Some of our results on zero- $G$ Hydrogen spectrum

- The Dirac Hamiltonian on zero- $G$  Kerr Newman spacetime is **essentially self-adjoint**
- The spectrum is **symmetric** about zero.
- Discrete spectrum is **non-empty** under some smallness conditions.
- The continuous spectrum is  $(-\infty, -m] \cup [m, \infty)$

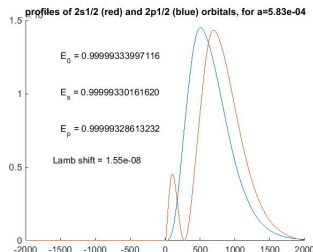
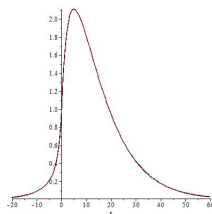
# Some of our results on zero- $G$ Hydrogen spectrum

- The Dirac Hamiltonian on zero- $G$  Kerr Newman spacetime is **essentially self-adjoint**
- The spectrum is **symmetric** about zero.
- Discrete spectrum is **non-empty** under some smallness conditions.
- The continuous spectrum is  $(-\infty, -m] \cup [m, \infty)$
- Profile of the positive energy ground state shows the tiger's tail!



# Some of our results on zero-G Hydrogen spectrum

- The Dirac Hamiltonian on zero-G Kerr Newman spacetime is **essentially self-adjoint**
- The spectrum is **symmetric** about zero.
- Discrete spectrum is **non-empty** under some smallness conditions.
- The continuous spectrum is  $(-\infty, -m] \cup [m, \infty)$
- Profile of the positive energy ground state shows the tiger's tail!



- Excited states. Numerical approximation. Hyperfine splitting and Lamb shift **without QED!**

# Back to Einstein, Rosen, and Weyl

- Recall: The ring singularity of the  $zGKN$  spacetime

# Back to Einstein, Rosen, and Weyl

- Recall: The ring singularity of the  $z$ GKN spacetime
- It connects the two sheets of the (otherwise vacuum) spacetime

# Back to Einstein, Rosen, and Weyl

- Recall: The ring singularity of the  $z$ GKN spacetime
- It connects the two sheets of the (otherwise vacuum) spacetime
- It's the locus of singularities for the metric, so the two-sheeted spacetime is defined outside a timelike 2-dim tube (circle  $\times$  real line)

# Back to Einstein, Rosen, and Weyl

- Recall: The ring singularity of the  $z$ GKN spacetime
- It connects the two sheets of the (otherwise vacuum) spacetime
- It's the locus of singularities for the metric, so the two-sheeted spacetime is defined outside a timelike 2-dim tube (circle  $\times$  real line)
- It's positively charged in one sheet, negatively charged in the other

# Back to Einstein, Rosen, and Weyl

- Recall: The ring singularity of the  $z$ GKN spacetime
- It connects the two sheets of the (otherwise vacuum) spacetime
- It's the locus of singularities for the metric, so the two-sheeted spacetime is defined outside a timelike 2-dim tube (circle  $\times$  real line)
- It's positively charged in one sheet, negatively charged in the other
- Our radically new idea: **Electron and Positron are not distinct particles but in fact “two different sides of the same coin”**

# Back to Einstein, Rosen, and Weyl

- Recall: The ring singularity of the  $z$ GKN spacetime
- It connects the two sheets of the (otherwise vacuum) spacetime
- It's the locus of singularities for the metric, so the two-sheeted spacetime is defined outside a timelike 2-dim tube (circle  $\times$  real line)
- It's positively charged in one sheet, negatively charged in the other
- Our radically new idea: **Electron and Positron are not distinct particles but in fact “two different sides of the same coin”**
- This resolves the paradox that Dirac's equation “for the electron” also seems to describe “a positron” in many situations, while it is a true one-particle equation.

# Topo-spin

- We introduce the notion of **TOPOLOGICAL SPIN** in analogy to Heisenberg's iso-spin.

# Topo-spin

- We introduce the notion of **TOPOLOGICAL SPIN** in analogy to Heisenberg's iso-spin.
- we identify the ring singularity of the zGKN spacetime with a **two-faced particle**, one that appears as an electron in one sheet and a positron in another sheet

# Topo-spin

- We introduce the notion of **TOPOLOGICAL SPIN** in analogy to Heisenberg's iso-spin.
- we identify the ring singularity of the zGKN spacetime with a **two-faced particle**, one that appears as an electron in one sheet and a positron in another sheet
- The radius of the ring = anomalous magnetic moment of the electron

# Topo-spin

- We introduce the notion of **TOPOLOGICAL SPIN** in analogy to Heisenberg's iso-spin.
- we identify the ring singularity of the zGKN spacetime with a **two-faced particle**, one that appears as an electron in one sheet and a positron in another sheet
- The radius of the ring = anomalous magnetic moment of the electron
- Can we formulate a quantum law of motion for the (center of the) ring?

# Topo-spin

- We introduce the notion of **TOPOLOGICAL SPIN** in analogy to Heisenberg's iso-spin.
- we identify the ring singularity of the zGKN spacetime with a **two-faced particle**, one that appears as an electron in one sheet and a positron in another sheet
- The radius of the ring = anomalous magnetic moment of the electron
- Can we formulate a quantum law of motion for the (center of the) ring?
- **YES! By relativity, it is the one that we have discussed!** (at least for quasi-static motions)

# Topo-spin

- We introduce the notion of **TOPOLOGICAL SPIN** in analogy to Heisenberg's iso-spin.
- we identify the ring singularity of the zGKN spacetime with a **two-faced particle**, one that appears as an electron in one sheet and a positron in another sheet
- The radius of the ring = anomalous magnetic moment of the electron
- Can we formulate a quantum law of motion for the (center of the) ring?
- YES! By relativity, it is the one that we have discussed! (at least for quasi-static motions)
- Anti-symmetry of the Dirac Hamiltonian with respect to topo-spin flips gives rise to the *matter-antimatter duality*

# Proof of symmetry of spectrum

- Let  $H$  be a matrix with a real eigenvalue  $E$  and a corresponding eigenvector  $\Psi$

$$H\Psi = E\Psi$$

# Proof of symmetry of spectrum

- Let  $H$  be a matrix with a real eigenvalue  $E$  and a corresponding eigenvector  $\Psi$

$$H\Psi = E\Psi$$

- Suppose there is another matrix  $C$  that anti-commutes with  $H$ , and  $C\Psi \neq 0$ :

$$CH + HC = 0$$

# Proof of symmetry of spectrum

- Let  $H$  be a matrix with a real eigenvalue  $E$  and a corresponding eigenvector  $\Psi$

$$H\Psi = E\Psi$$

- Suppose there is another matrix  $C$  that anti-commutes with  $H$ , and  $C\Psi \neq 0$ :

$$CH + HC = 0$$

- Then  $-E$  is also an eigenvalue of  $H$ , with eigenvector  $C\Psi$ .

# Proof of symmetry of spectrum

- Let  $H$  be a matrix with a real eigenvalue  $E$  and a corresponding eigenvector  $\Psi$

$$H\Psi = E\Psi$$

- Suppose there is another matrix  $C$  that anti-commutes with  $H$ , and  $C\Psi \neq 0$ :

$$CH + HC = 0$$

- Then  $-E$  is also an eigenvalue of  $H$ , with eigenvector  $C\Psi$ .
- We found a  $C$  that does the job for our Hamiltonian  $H$ .

# Proof of symmetry of spectrum

- Let  $H$  be a matrix with a real eigenvalue  $E$  and a corresponding eigenvector  $\Psi$

$$H\Psi = E\Psi$$

- Suppose there is another matrix  $C$  that anti-commutes with  $H$ , and  $C\Psi \neq 0$ :

$$CH + HC = 0$$

- Then  $-E$  is also an eigenvalue of  $H$ , with eigenvector  $C\Psi$ .
- We found a  $C$  that does the job for our Hamiltonian  $H$ .
- $C$  is a topo-spin flip!

# Proof of symmetry of spectrum

- Let  $H$  be a matrix with a real eigenvalue  $E$  and a corresponding eigenvector  $\Psi$

$$H\Psi = E\Psi$$

- Suppose there is another matrix  $C$  that anti-commutes with  $H$ , and  $C\Psi \neq 0$ :

$$CH + HC = 0$$

- Then  $-E$  is also an eigenvalue of  $H$ , with eigenvector  $C\Psi$ .
- We found a  $C$  that does the job for our Hamiltonian  $H$ .
- $C$  is a topo-spin flip!
- Eigenfunctions with positive energy are 99% supported in one sheet, and those with negative energy are 99% supported in the other sheet

# Summary

- Zero- $G$  general relativity is NOT necessarily special relativity, and zero- $G$  spacetimes NOT necessarily merely “wavy” perturbations of Minkowski spacetime.

# Summary

- Zero- $G$  general relativity is NOT necessarily special relativity, and zero- $G$  spacetimes NOT necessarily merely “wavy” perturbations of Minkowski spacetime.
- Topologically non-trivial spacetimes should be taken seriously

# Summary

- Zero- $G$  general relativity is NOT necessarily special relativity, and zero- $G$  spacetimes NOT necessarily merely “wavy” perturbations of Minkowski spacetime.
- Topologically non-trivial spacetimes should be taken seriously
- The Dirac equation on  $z$ GKN is well-posed; **Naked singularity means no harm!**

# Summary

- Zero- $G$  general relativity is NOT necessarily special relativity, and zero- $G$  spacetimes NOT necessarily merely “wavy” perturbations of Minkowski spacetime.
- Topologically non-trivial spacetimes should be taken seriously
- The Dirac equation on zGKN is well-posed; **Naked singularity means no harm!**
- The Dirac Hamiltonian on zGKN has **symmetric spectrum** with **scattering and bound states**

# Summary

- Zero- $G$  general relativity is NOT necessarily special relativity, and zero- $G$  spacetimes NOT necessarily merely “wavy” perturbations of Minkowski spacetime.
- Topologically non-trivial spacetimes should be taken seriously
- The Dirac equation on zGKN is well-posed; **Naked singularity means no harm!**
- The Dirac Hamiltonian on zGKN has **symmetric spectrum** with **scattering and bound states**
- Novel proposal: **Dirac's equation describes a single “particle / anti-particle” structure: two “topo-spin” states**

# Outlook

- Characterize the discrete Dirac spectrum on  $zGKN$  completely

# Outlook

- Characterize the discrete Dirac spectrum on  $zGKN$  completely
- Do lots of numerical experiments

# Outlook

- Characterize the discrete Dirac spectrum on  $zGKN$  completely
- Do lots of numerical experiments
- Study zero- $G$  general-relativistic Dirac spectrum for “Positronium” (2 rings, 4 sheets?)

# Outlook

- Characterize the discrete Dirac spectrum on  $zGKN$  completely
- Do lots of numerical experiments
- Study zero- $G$  general-relativistic Dirac spectrum for “Positronium” (2 rings, 4 sheets?)
- Limit of  $\infty$  many rings may exhibit “ferro-topological phase transition”: **symmetry breaking** (Explanation of broken particle / anti-particle symmetry in our world?).

# Outlook

- Characterize the discrete Dirac spectrum on zGKN completely
- Do lots of numerical experiments
- Study zero- $G$  general-relativistic Dirac spectrum for “Positronium” (2 rings, 4 sheets?)
- Limit of  $\infty$  many rings may exhibit “ferro-topological phase transition”: **symmetry breaking** (Explanation of broken particle / anti-particle symmetry in our world?).
- Turn gravity back on, within a nonlinear EM theory. (Perturbation method?)

# Fin!

THANK YOU FOR LISTENING!

# Proof of existence of discrete spectrum

$$\bullet \Psi(t, r, \theta, \varphi) = R(r)S(\theta)e^{-i(Et - \kappa\varphi)} \begin{pmatrix} \cos(\Theta(\theta)/2)e^{+i\Omega(r)/2} \\ \sin(\Theta(\theta)/2)e^{-i\Omega(r)/2} \\ \cos(\Theta(\theta)/2)e^{-i\Omega(r)/2} \\ \sin(\Theta(\theta)/2)e^{+i\Omega(r)/2} \end{pmatrix}$$

# Proof of existence of discrete spectrum

- $$\Psi(t, r, \theta, \varphi) = R(r)S(\theta)e^{-i(Et - \kappa\varphi)} \begin{pmatrix} \cos(\Theta(\theta)/2)e^{+i\Omega(r)/2} \\ \sin(\Theta(\theta)/2)e^{-i\Omega(r)/2} \\ \cos(\Theta(\theta)/2)e^{-i\Omega(r)/2} \\ \sin(\Theta(\theta)/2)e^{+i\Omega(r)/2} \end{pmatrix}$$
- $$\begin{cases} d\Omega/dr = 2\frac{mr}{\Delta} \cos \Omega + 2\frac{\lambda}{\Delta} \sin \Omega + 2\frac{a\kappa + \gamma r}{\Delta^2} - 2E \\ d(\ln R)/dr = \frac{mr}{\Delta} \sin \Omega - \frac{\lambda}{\Delta} \cos \Omega \\ \begin{cases} d\Theta/d\theta = 2(\lambda - m a \cos \theta \cos \Theta + (aE \sin \theta - \frac{\kappa}{\sin \theta}) \sin \Theta) \\ d(\ln S)/d\theta = -m a \cos \theta \sin \Theta - (aE \sin \theta - \frac{\kappa}{\sin \theta}) \cos \Theta. \end{cases} \end{cases}$$

# Proof of existence of discrete spectrum

- $\Psi(t, r, \theta, \varphi) = R(r)S(\theta)e^{-i(Et - \kappa\varphi)} \begin{pmatrix} \cos(\Theta(\theta)/2)e^{+i\Omega(r)/2} \\ \sin(\Theta(\theta)/2)e^{-i\Omega(r)/2} \\ \cos(\Theta(\theta)/2)e^{-i\Omega(r)/2} \\ \sin(\Theta(\theta)/2)e^{+i\Omega(r)/2} \end{pmatrix}$
- $$\begin{cases} d\Omega/dr = 2\frac{mr}{\Delta} \cos \Omega + 2\frac{\lambda}{\Delta} \sin \Omega + 2\frac{a\kappa + \gamma r}{\Delta^2} - 2E \\ d(\ln R)/dr = \frac{mr}{\Delta} \sin \Omega - \frac{\lambda}{\Delta} \cos \Omega \\ \begin{cases} d\Theta/d\theta = 2(\lambda - ma \cos \theta \cos \Theta + (aE \sin \theta - \frac{\kappa}{\sin \theta}) \sin \Theta) \\ d(\ln S)/d\theta = -ma \cos \theta \sin \Theta - (aE \sin \theta - \frac{\kappa}{\sin \theta}) \cos \Theta. \end{cases} \end{cases}$$
- $\Psi \in L^2$  iff:
 
$$\begin{cases} \Omega(-\infty) = -\pi + \cos^{-1}(E), & \Omega(\infty) = -\cos^{-1}(E) \\ \Theta(0) = 0, & \Theta(\pi) = -\pi. \end{cases}$$

# Flows on a finite cylinder

$$\bullet \begin{cases} \dot{\theta} = \sin \theta \\ \dot{\Theta} = -2a \sin \theta \cos \theta \cos \Theta + 2aE \sin^2 \theta \sin \Theta - 2\kappa \sin \Theta \\ \quad + 2\lambda \sin \theta \\ \dot{\xi} = \cos^2 \xi \\ \dot{\Omega} = 2a \sin \xi \cos \Omega + 2\lambda \cos \xi \sin \Omega + 2\gamma \sin \xi \cos \xi \\ \quad + 2\kappa \cos^2 \xi - 2aE \end{cases}$$

# Flows on a finite cylinder

- $$\begin{cases} \dot{\theta} &= \sin \theta \\ \dot{\Theta} &= -2a \sin \theta \cos \theta \cos \Theta + 2aE \sin^2 \theta \sin \Theta - 2\kappa \sin \Theta \\ &\quad + 2\lambda \sin \theta \end{cases}$$
- $$\begin{cases} \dot{\xi} &= \cos^2 \xi \\ \dot{\Omega} &= 2a \sin \xi \cos \Omega + 2\lambda \cos \xi \sin \Omega + 2\gamma \sin \xi \cos \xi \\ &\quad + 2\kappa \cos^2 \xi - 2aE \end{cases}$$
- Parameter-dependent flow on a cylinder

$$\begin{cases} \dot{x} &= f(x) \\ \dot{y} &= g_{\mu}(x, y) \end{cases} \quad (x, y) \in [x_-, x_+] \times \mathbb{S}^1$$

# Flows on a finite cylinder

- $$\begin{cases} \dot{\theta} &= \sin \theta \\ \dot{\Theta} &= -2a \sin \theta \cos \theta \cos \Theta + 2aE \sin^2 \theta \sin \Theta - 2\kappa \sin \Theta \\ &\quad + 2\lambda \sin \theta \end{cases}$$
- $$\begin{cases} \dot{\xi} &= \cos^2 \xi \\ \dot{\Omega} &= 2a \sin \xi \cos \Omega + 2\lambda \cos \xi \sin \Omega + 2\gamma \sin \xi \cos \xi \\ &\quad + 2\kappa \cos^2 \xi - 2aE \end{cases}$$
- Parameter-dependent flow on a cylinder

$$\begin{cases} \dot{x} &= f(x) \\ \dot{y} &= g_{\mu}(x, y) \end{cases} \quad (x, y) \in [x_-, x_+] \times \mathbb{S}^1$$

- Two equilibrium points on each boundary:  $f(x_-) = f(x_+) = 0$  and  $g_{\mu}(x_{\pm}, y) = 0 \implies y \in \{s_{\pm}, n_{\pm}\}$

# Looking for Heteroclinic Orbits

- Nodes  $N_{\pm} = (x_{\pm}, n_{\pm})$  and Saddles  $S_{\pm} = (x_{\pm}, s_{\pm})$

# Looking for Heteroclinic Orbits

- Nodes  $N_{\pm} = (x_{\pm}, n_{\pm})$  and Saddles  $S_{\pm} = (x_{\pm}, s_{\pm})$
- Existence of  $N_-$  to  $S_+$  and  $S_-$  to  $N_+$  connecting orbits (stable/unstable and center manifold theory.)

# Looking for Heteroclinic Orbits

- Nodes  $N_{\pm} = (x_{\pm}, n_{\pm})$  and Saddles  $S_{\pm} = (x_{\pm}, s_{\pm})$
- Existence of  $N_-$  to  $S_+$  and  $S_-$  to  $N_+$  connecting orbits (stable/unstable and center manifold theory.)
- Existence of  $L^2$  eigenfunction iff there is an orbit connecting the two saddles:  $S_-$  to  $S_+$ .

# Looking for Heteroclinic Orbits

- Nodes  $N_{\pm} = (x_{\pm}, n_{\pm})$  and Saddles  $S_{\pm} = (x_{\pm}, s_{\pm})$
- Existence of  $N_-$  to  $S_+$  and  $S_-$  to  $N_+$  connecting orbits (stable/unstable and center manifold theory.)
- Existence of  $L^2$  eigenfunction iff there is an orbit connecting the two saddles:  $S_-$  to  $S_+$ .
- The Corridor formed by the two SN connectors.



# Looking for Heteroclinic Orbits

- Nodes  $N_{\pm} = (x_{\pm}, n_{\pm})$  and Saddles  $S_{\pm} = (x_{\pm}, s_{\pm})$
- Existence of  $N_-$  to  $S_+$  and  $S_-$  to  $N_+$  connecting orbits (stable/unstable and center manifold theory.)
- Existence of  $L^2$  eigenfunction iff there is an orbit connecting the two saddles:  $S_-$  to  $S_+$ .
- The Corridor formed by the two SN connectors.



- Saddle-Saddle connector exists iff the corridor collapses.

# Topological Methods in Dynamical Systems

- $\mu$ -dependent flow on a cylinder

# Topological Methods in Dynamical Systems

- $\mu$ -dependent flow on a cylinder
- $\mathbf{a}(\mu)$  = area of corridor

# Topological Methods in Dynamical Systems

- $\mu$ -dependent flow on a cylinder
- $\mathbf{a}(\mu)$  = area of corridor
- $\mathbf{w}(\mu)$  = winding number of corridor

# Topological Methods in Dynamical Systems

- $\mu$ -dependent flow on a cylinder
- $\mathbf{a}(\mu)$  = area of corridor
- $\mathbf{w}(\mu)$  = winding number of corridor
- $\mathbf{a}$  is a continuous function of  $\mu$

# Topological Methods in Dynamical Systems

- $\mu$ -dependent flow on a cylinder
- $\mathbf{a}(\mu)$  = area of corridor
- $\mathbf{w}(\mu)$  = winding number of corridor
- $\mathbf{a}$  is a continuous function of  $\mu$
- $\mathbf{a} > 0$  iff  $\mathbf{w} \geq 1$  (Green's theorem).

# Topological Methods in Dynamical Systems

- $\mu$ -dependent flow on a cylinder
- $\mathbf{a}(\mu)$  = area of corridor
- $\mathbf{w}(\mu)$  = winding number of corridor
- $\mathbf{a}$  is a continuous function of  $\mu$
- $\mathbf{a} > 0$  iff  $\mathbf{w} \geq 1$  (Green's theorem).
- $\mathbf{a} < 0$  iff  $\mathbf{w} \leq 0$ .

# Topological Methods in Dynamical Systems

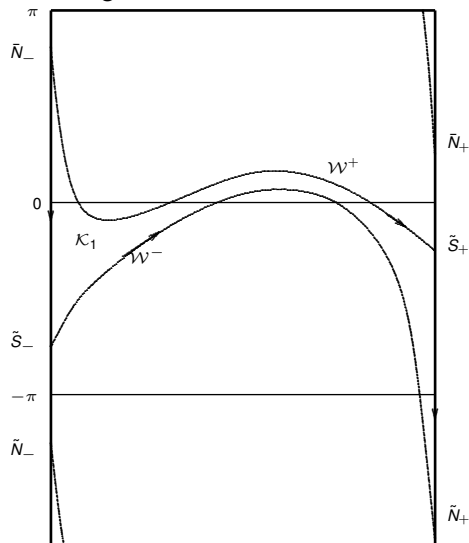
- $\mu$ -dependent flow on a cylinder
- $\mathbf{a}(\mu)$  = area of corridor
- $\mathbf{w}(\mu)$  = winding number of corridor
- $\mathbf{a}$  is a continuous function of  $\mu$
- $\mathbf{a} > 0$  iff  $\mathbf{w} \geq 1$  (Green's theorem).
- $\mathbf{a} < 0$  iff  $\mathbf{w} \leq 0$ .
- $\mathbf{a} = 0$  iff corridor is empty (i.e. there is a saddle connection.)

# Topological Methods in Dynamical Systems

- $\mu$ -dependent flow on a cylinder
- $\mathbf{a}(\mu)$  = area of corridor
- $\mathbf{w}(\mu)$  = winding number of corridor
- $\mathbf{a}$  is a continuous function of  $\mu$
- $\mathbf{a} > 0$  iff  $\mathbf{w} \geq 1$  (Green's theorem).
- $\mathbf{a} < 0$  iff  $\mathbf{w} \leq 0$ .
- $\mathbf{a} = 0$  iff corridor is empty (i.e. there is a saddle connection.)
- Construction of barriers to prove existence of corridors with given winding number.

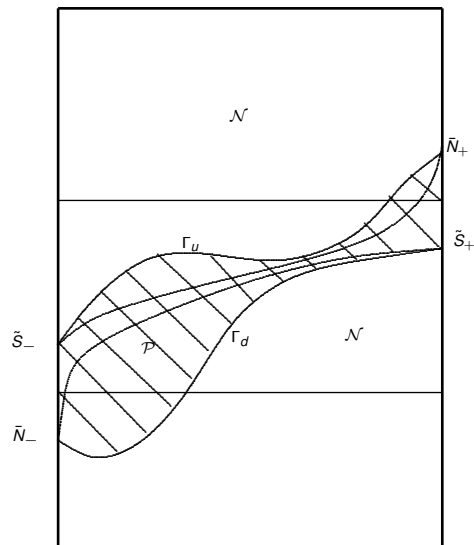
# Area and Winding Number for Corridors

Working in the universal cover of the cylinder:

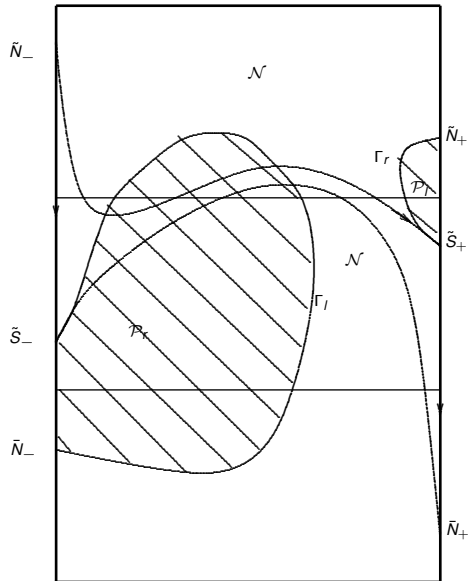


# Topology of Nullclines

Orbits must increase while in the shaded region



# Change in Nullcline Topology and Corridor Winding #



# Barrier construction

